

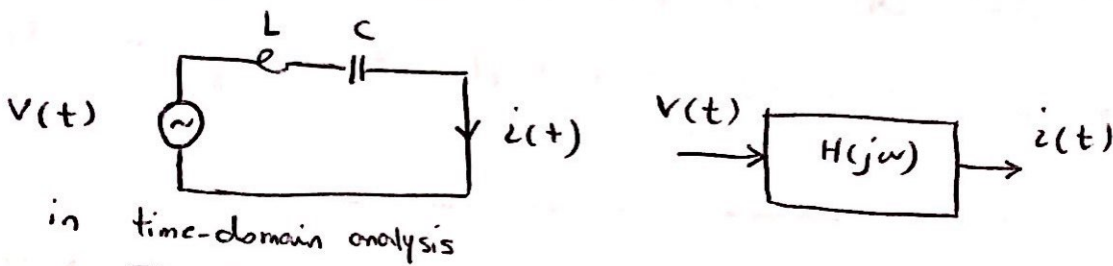
① What is resonant? $\left\{ \begin{array}{l} \text{Amp. over time} \\ \text{Freq. dependence} \end{array} \right.$

② Example of resonant system (series LC)

↳ How to find if it is resonant? Excite @ all freqs and see which one is selected.

↳ Excite @ oscillating freq. and see amplification.

↳ The transfer function of the resonant system.

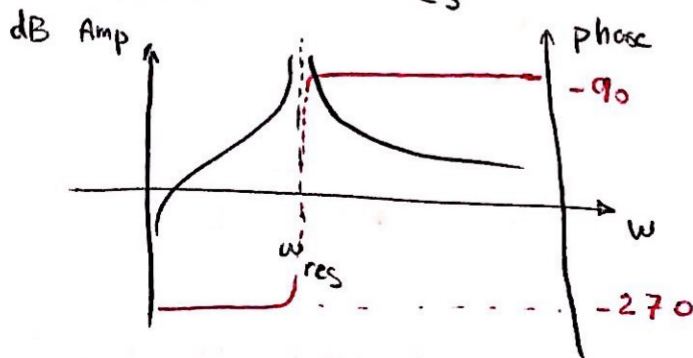


• we saw upon excitation at resonance (i.e., $\omega_{res} = \frac{1}{\sqrt{LC}}$) the oscillating response grows with time:

(i.e., when $v(t) = \sin(\omega_{res} t) u(t)$, then $i(t) = \frac{L}{2} t \sin(\omega_{res} t) u(t)$)

• Now let's find the transfer function of this resonant system:

$$H(s) = \frac{I(s)}{V(s)} = \frac{1}{Ls + \frac{1}{Cs}} \rightarrow H(j\omega) = \frac{1}{j\omega L + \frac{1}{j\omega C}} = \frac{-j\omega/L}{\omega^2 - \frac{1}{LC}}$$



* Resonance is a peak @ ω_{res} !

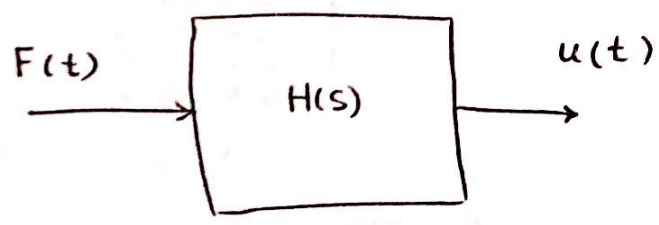
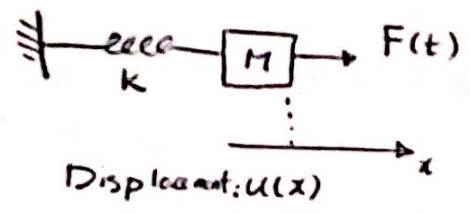
* The phase changes by 180° @ ω_{res} !

* Generalization: Any system with $H(s) = \frac{N(s)}{D(s)}$, where $D(s)$ has a complex-conjugate pole pair of $\pm j\omega_0$, is resonant, with a resonance freq. of ω_0 .

HW2, p1: Prove this statement by time-domain analysis. (Hint: excite the system @ ω_0 , and show the amplification happens!)

↳ Immediate extension: if $H(s)$ has K complex-conjugate poles at $\pm j\omega_i$ ($i=1$ to K), then the system has K resonance modes!

Example 2: Mechanical resonant system (harmonic oscillator):



Newton's 2nd law

$$M \frac{d^2 u}{dt^2} + \underbrace{Ku}_{\text{Hooke's law}} = F \xrightarrow{\mathcal{L}} U(s) (Ms^2 + K) = F(s) \rightarrow H(s) = \frac{U(s)}{F(s)} = \frac{1}{Ms^2 + K}$$

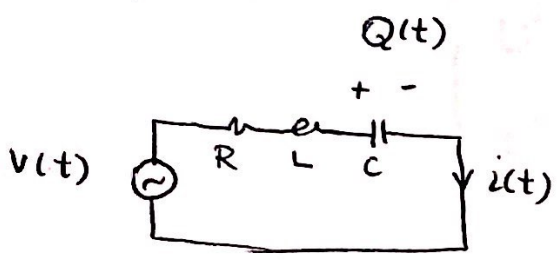
→ poles: $s = \pm j \sqrt{\frac{K}{M}}$ → A resonant system; $\omega_{res} = \sqrt{\frac{K}{M}}$.

HW2, p2: We want to make an actuator with largest possible displacement @ lowest input force! what is your solution?

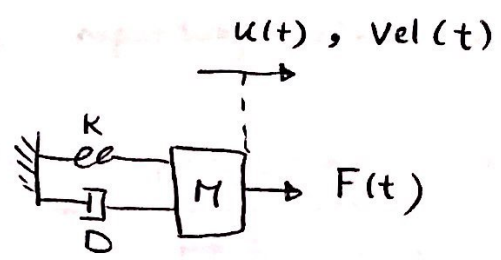
Damped oscillation / resonance: In reality, all resonant systems have at least one damping mechanism that attenuates the signal.

- ↳ Electrical: resistance, elec-mag radiation, etc.
- ↳ Mechanical: air damping, thermal noise, friction, etc.

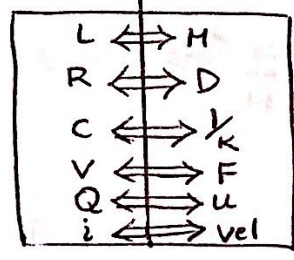
Example 3:



$$L \frac{d^2 Q}{dt^2} + R \frac{dQ}{dt} + \frac{1}{C} Q = V(t)$$



$$M \frac{d^2 u}{dt^2} + D \frac{du}{dt} + K u = F(t)$$



* Resonant systems are "dynamically equivalent"!

↳ All the math interchangeable!

↳ Note: loss/attenuation in both are essentially same { R & D }

* Now time & freq. analysis in presence of loss/damping.

↳ Assuming impulse input & finding u(t) {or equivalently Q(t)}.

$$\mathcal{L} \rightarrow \begin{cases} (L s^2 + R s + \frac{1}{C}) Q(s) = 1 \rightarrow Q(s) = \frac{1}{L s^2 + R s + \frac{1}{C}} \\ (M s^2 + D s + K) U(s) = 1 \rightarrow U(s) = \frac{1}{M s^2 + D s + K} \end{cases}$$

* Poles of $V(s)$:
$$\frac{-D \pm \sqrt{D^2 - 4MK}}{2M}$$

↳ As we know from circuits / ODE: three scenarios could happen:

$$\left\{ \begin{array}{l} D^2 > 4MK \rightarrow \text{two real-roots} \rightarrow \text{over-damped.} \\ D^2 = 4MK \rightarrow \text{repeating real-roots} \rightarrow \text{critically damped.} \\ D^2 < 4MK \rightarrow \text{two complex-conjugate roots! underdamped.} \end{array} \right.$$

↳ This is the required criteria for resonance system!

* Assuming under-damped case:

↳ poles are
$$-\frac{D}{2M} \pm j \left(\frac{K}{M} - \frac{D^2}{4M^2} \right)^{0.5}$$

Damping ratio.

* Defining:
$$\left\{ \begin{array}{l} \zeta = \frac{D}{2\sqrt{KM}} \\ \omega_0 = \sqrt{\frac{K}{M}} \end{array} \right. \rightarrow \text{poles:}$$

$$-\zeta\omega_0 \pm j\omega_0\sqrt{1-\zeta^2}$$

$$\omega_d = \omega_0\sqrt{1-\zeta^2}$$

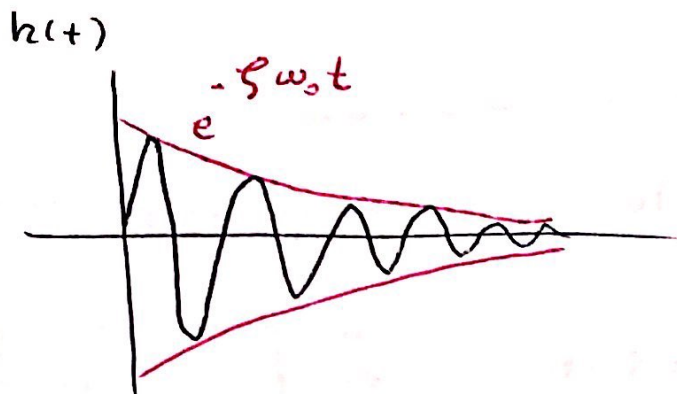
Damped res. freq.

(undamped) Natural resonance freq of the system

$$\Rightarrow u(t) = \sqrt{\frac{1}{KM}} e^{-\zeta\omega_0 t} \sin(\omega_d t + \varphi)$$

key take-aways:
$$\left\{ \begin{array}{l} 1) \text{ Damping: exponentially attenuates resonance oscillation} \\ 2) \text{ Damping: reduces resonance freq to } \omega_d! \end{array} \right.$$

* plotting impulse response: $u(t) = \frac{1}{\sqrt{KM}} e^{-\zeta\omega_0 t} \sin(\omega_d t + \varphi)$:



* Quality factor: The lower the damping rate, the higher the quality of resonator!

$$Q \triangleq 2\pi \frac{\text{Energy stored in the system}}{\text{Energy lost per cycle}}$$

(The quality factor)

* For the harmonic oscillator, energy stored in the system at the beginning of each cycle \downarrow is: $\frac{1}{2} K \left\{ \frac{1}{\sqrt{KM}} e^{-\zeta\omega_d NT} \right\}^2$
i.e. NT

HW2, P3: why the above statement is correct?

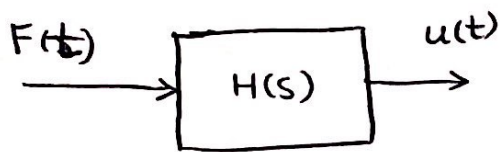
$$\Rightarrow Q = \frac{\frac{1}{2M} e^{-2\zeta\omega_d NT} \cdot 2\pi}{\frac{1}{2M} e^{-2\zeta\omega_d NT} - \frac{1}{2M} e^{-2\zeta\omega_d (N+1)T}} = \frac{2\pi}{1 - e^{-2\zeta\omega_d T}} \approx \frac{2\pi}{2\zeta\omega_d T}$$

$$\Rightarrow Q = \frac{2\pi}{2\zeta \underbrace{\omega_d T}_{=2\pi}} = \frac{1}{2\zeta}$$

$$Q = \frac{1}{2\zeta}$$

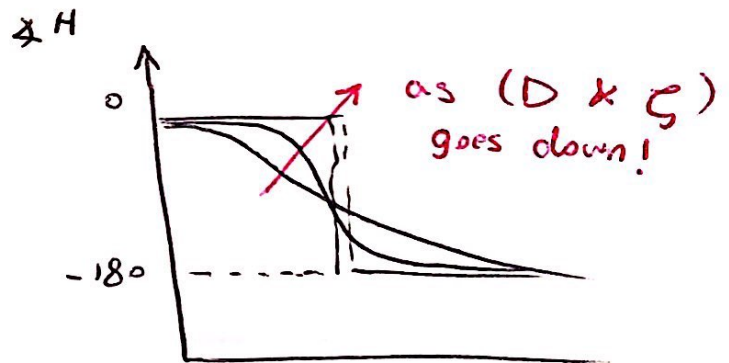
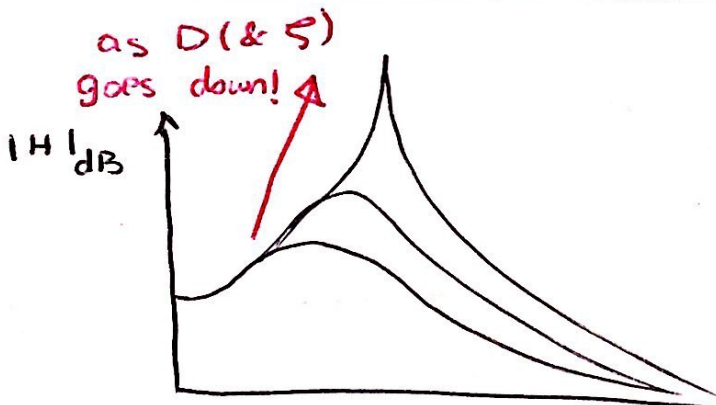
HW2, p 31: Derive the time-domain wave-form of damped resonant system for $\sin(\omega_d t)$ excitation!

* Now freq-domain analysis of damped-resonant system:



$$H(s) = \frac{1}{Ms + R + \frac{1}{Cs}}$$

$$\Rightarrow H(j\omega) = \frac{-\frac{1}{M}}{(\omega^2 - \omega_0^2) - j(2\zeta\omega_0\omega)}$$



the lower D & ζ

180°

* From freq. response, the sharper the peak & phase transition! [we can use the sharpness of the peak as a measure of resonator quality!]

Sharpness of the peak: $\frac{\omega_{res}}{\Delta\omega_{-3dB}}$; $\Delta\omega_{-3dB} = \omega_{-3dB,H} - \omega_{-3dB,L}$

* Let's find/calculate $\frac{\omega_{res}}{\Delta\omega_{-3dB}}$:

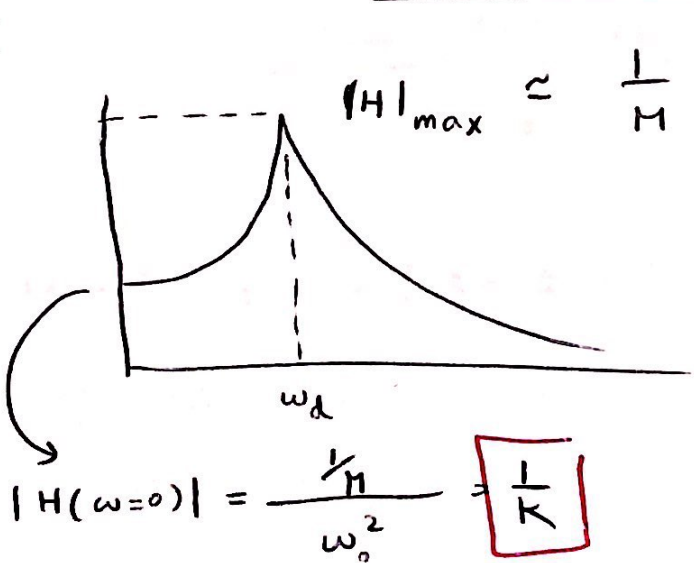
7

② ω_{-3dB} : $|H(j\omega)_{-3dB}| = \frac{1}{\sqrt{2}} |H(j\omega_d)| \approx \frac{1}{\sqrt{2}} \frac{1/M}{\omega_0^2 \zeta^2 - j2\zeta\omega_0\omega_d}$

$\Rightarrow \frac{|(\omega_{-3dB}^2 - \omega_0^2) - j(2\zeta\omega_0\omega_{-3dB})|}{\sqrt{2}} = \frac{1}{\omega_0^2 \zeta^2 - j2\zeta\omega_0\omega_d}$

HW2, Bonus (30 points): solve the eq. above to find $\omega_{-3dB,H}$ & $\omega_{-3dB,L}$. show that

$$\frac{\omega_d}{\Delta\omega_{-3dB}} = \frac{\omega_d}{\omega_{-3dB,H} - \omega_{-3dB,L}} = Q$$



$|H|_{max} \approx \frac{1}{M} \cdot \frac{1}{2\zeta\omega_0^2} = \frac{1}{2\zeta \cdot K} = \boxed{\frac{Q}{K}}$

* At resonance amp of displacement is Q times that of DC!

Conclusion of Freq. domain analysis :

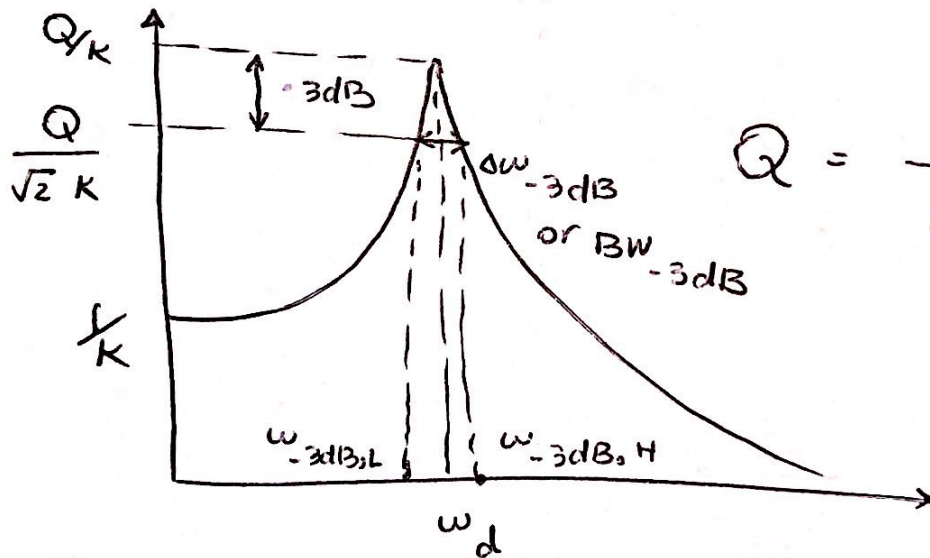
* The resonance shows as a peak @ $\omega_d = \omega_0 (1 - \zeta^2)^{0.5}$

* The sharpness of the peak is $Q = \frac{\omega_d}{\Delta \omega_{-3dB}} = \frac{1}{2\zeta}$

* For $\frac{U(j\omega)}{F(j\omega)}$, the oscillation amplitude at ω_{res}

is Q times displacement @ $\omega=0$!

$$H(j\omega)_{dB} = \left(\frac{U(j\omega)}{F(j\omega)} \right)_{dB}$$



$$Q = \frac{\omega_d}{\Delta \omega_{-3dB}}$$

H W 2, Bonus (30 points) : Prove Q can be extracted from $H(j\omega)$ phase slope at ω_{res} :

$$\text{i.e., } Q = \frac{\omega_{res}}{2} \cdot \frac{d\angle H(j\omega)}{d\omega}$$